

# Mixed Mode Oscillations Underlie Bursting in Pituitary Cells

Richard Bertram

Department of Mathematics  
Florida State University  
Tallahassee, Florida

# Collaborators and Support

Wondimu Teka

Joël Tabak



Florida State University

Theodore Vo

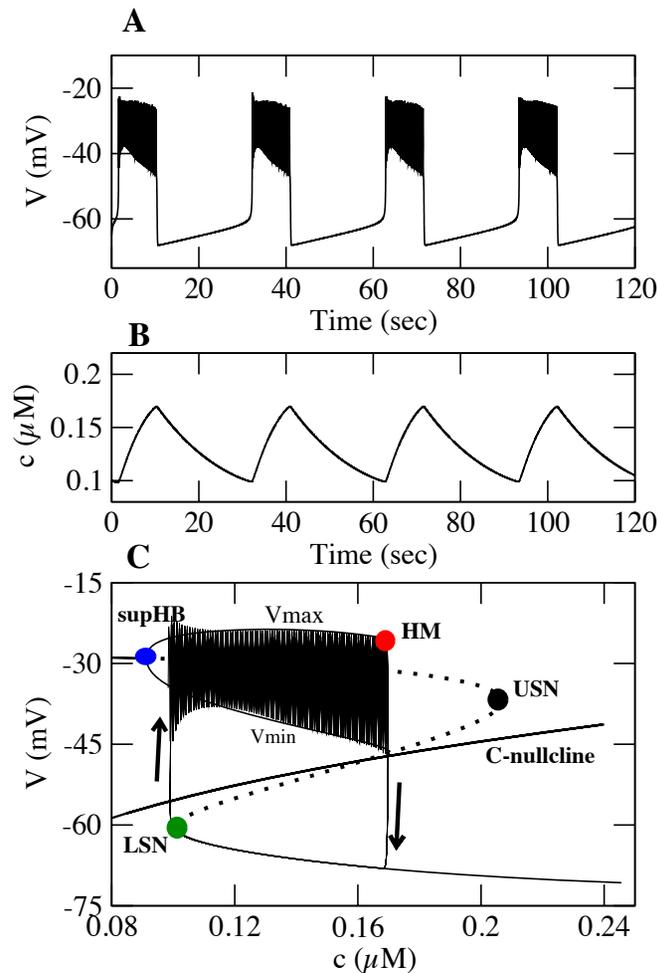
Martin Wechselberger



University of Sydney

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# Fast-Slow Analysis: a Powerful Tool for Understanding Plateau Bursting



$$\dot{V} = f(V, n, c)$$

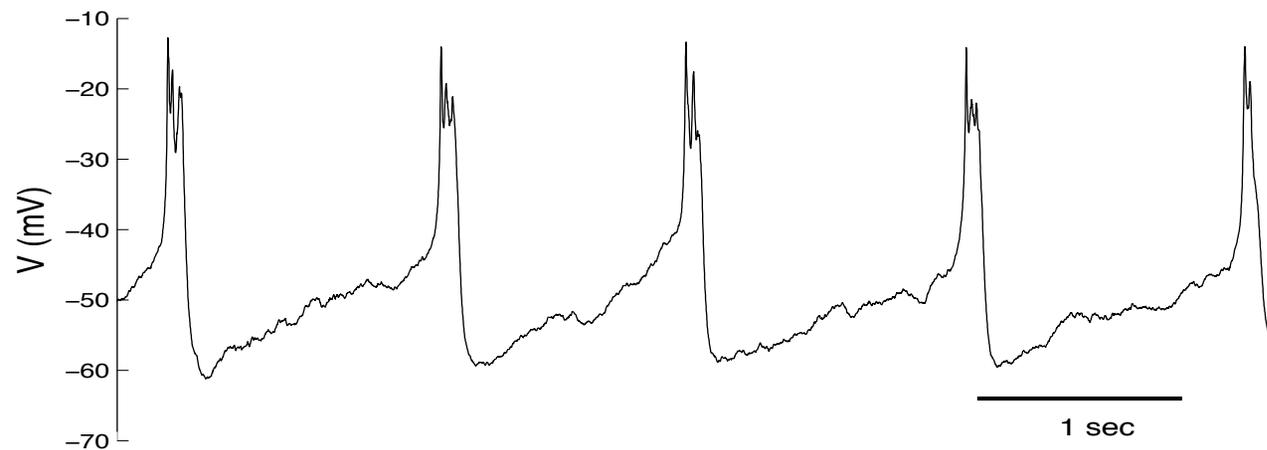
$$\dot{n} = g(V, n)$$

$$\dot{c} = \varepsilon h(V, c)$$

Analysis in the limit

$$\varepsilon \rightarrow 0$$

# Pseudo-Plateau Bursting Occurs in Some Pituitary Cells

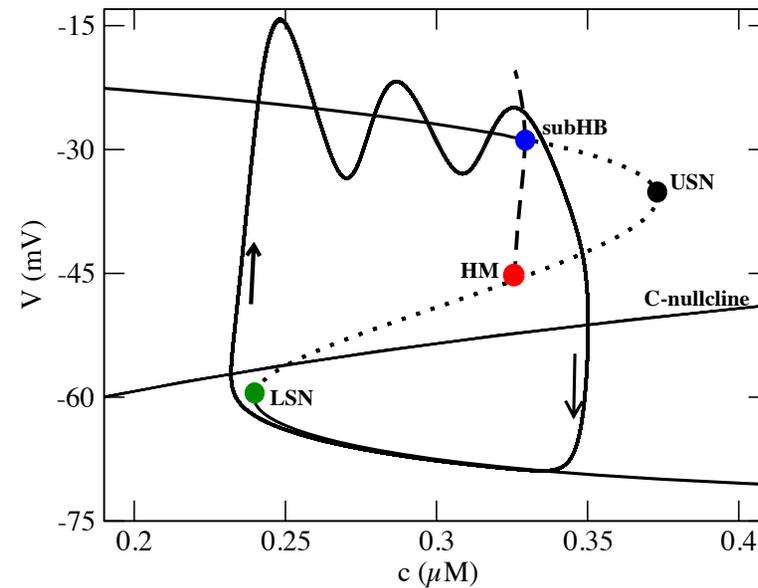
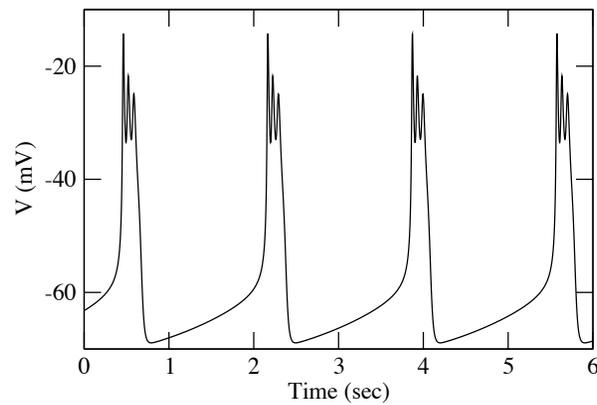


Electrical recording from a GH4 pituitary cell line

Bursts are short and the spikes have very small amplitude.

Characteristic of bursting in pituitary lactotrophs and somatotrophs

# Pseudo-Plateau Bursting Presents New Challenges



Trajectory does not follow the z-curve, and  
there is no periodic spiking branch!

# An Alternate Approach

$$\varepsilon \dot{V} = f(V, n, c)$$

$$\dot{n} = g(V, n)$$

$$\dot{c} = h(V, c)$$

Analyze the **reduced system** obtained in the limit  $\varepsilon \rightarrow 0$

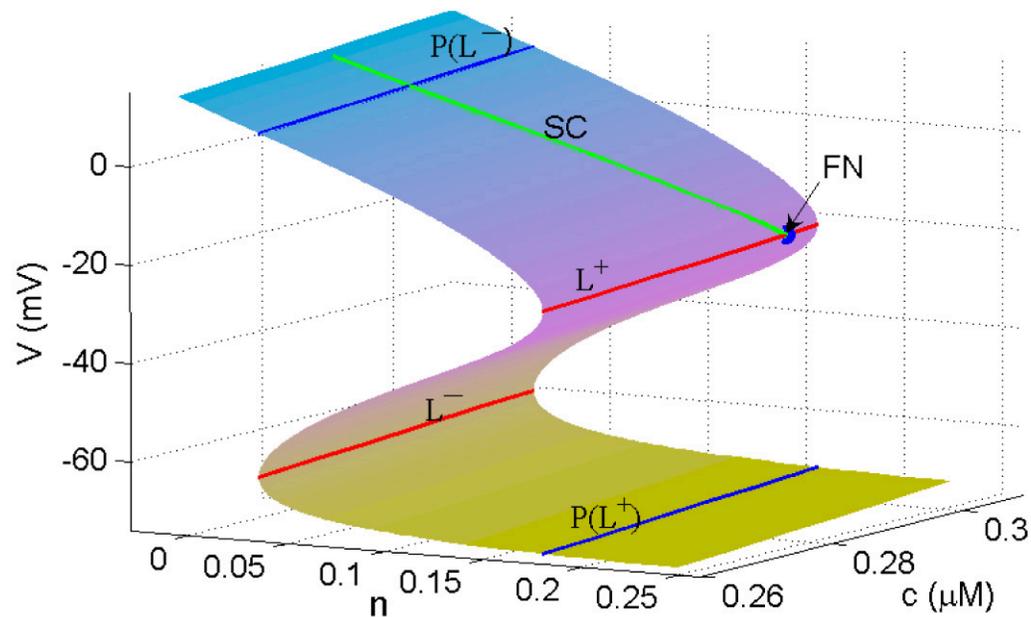
Voltage  $V$  is in a state of **quasi-equilibrium** with  $n$  and  $c$

# The Critical Manifold

Surface in 3-space where  $V$  is at quasi-equilibrium

RHS of V-ODE:  $f(V, n, c) \equiv -(I_{Ca} + I_K + I_{SK} + I_{BK})$

Critical manifold:  $S \equiv \{(V, n, c) \in \mathfrak{R}^3 : f(V, n, c) = 0\}$



# The Reduced and Desingularized Systems

RHS of V-ODE:  $f(V, n, c) \equiv -(I_{Ca} + I_K + I_{SK} + I_{BK})$

Critical manifold:  $S \equiv \{(V, n, c) \in \mathfrak{R}^3 : f(V, n, c) = 0\}$

Dynamics on S:  $\frac{d}{dt} f(V, n, c) = \frac{d}{dt} 0$

Reduced system:  $-\frac{\partial f}{\partial V} \frac{dV}{dt} = g(V, n) \frac{\partial f}{\partial n} + h(V, c) \frac{\partial f}{\partial c}$   
 $\frac{dc}{dt} = h(V, c)$

*(Note: A red arrow points from the text "=0 on folds" to the  $-\frac{\partial f}{\partial V} \frac{dV}{dt}$  term. Two blue arrows point from the  $\frac{dn}{dt}$  and  $\frac{dc}{dt}$  terms in the dynamics equation above to the  $\frac{\partial f}{\partial n}$  and  $\frac{\partial f}{\partial c}$  terms respectively.)*

=0 on folds

With  $n$  satisfying

$$f(V, n, c) = 0$$

Desingularized system:  $\frac{dV}{d\tau} = g(V, n) \frac{\partial f}{\partial n} + h(V, c) \frac{\partial f}{\partial c} \equiv F(V, n, c)$

$$\frac{dc}{d\tau} = -h(V, c) \frac{\partial f}{\partial V}$$

with  $\tau \equiv -\left(\frac{\partial f}{\partial V}\right)^{-1} t$

# Equilibria of the Desingularized System

Equilibrium  
Conditions

$$\frac{dV}{d\tau} = F(V, n, c) = 0 \quad \frac{dc}{d\tau} = -h(V, c) \frac{\partial f}{\partial V} = 0$$

Regular  
Singularity

$$g(V, n) = 0$$
$$h(V, c) = 0$$

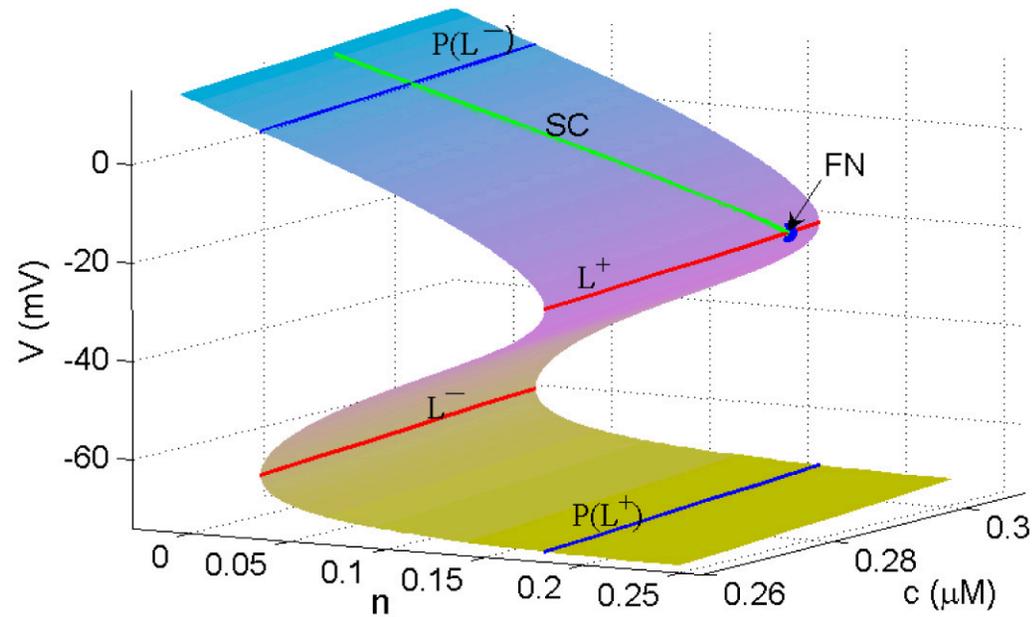
Folded  
Singularity

$$\frac{\partial f}{\partial V} = 0$$

(On a fold curve of S)

**Folded Node:** A folded singularity with two real eigenvalues

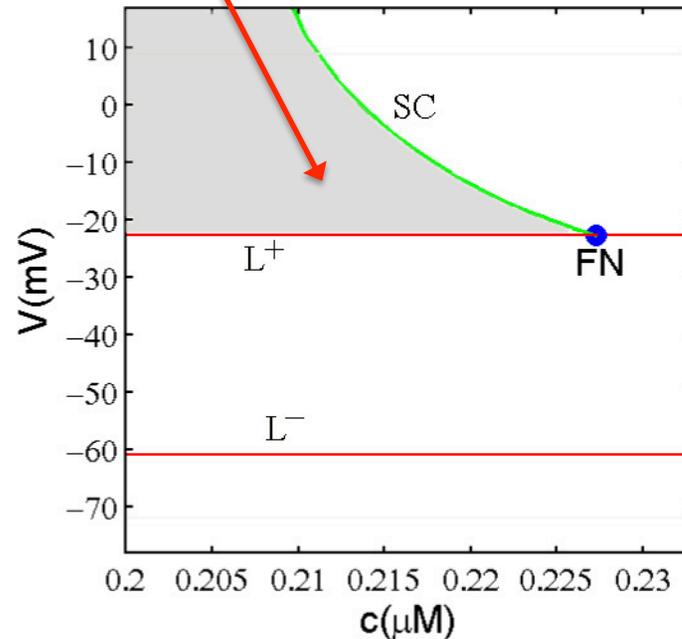
# There is a Folded Node on the Top Fold Curve



# Singular Canards

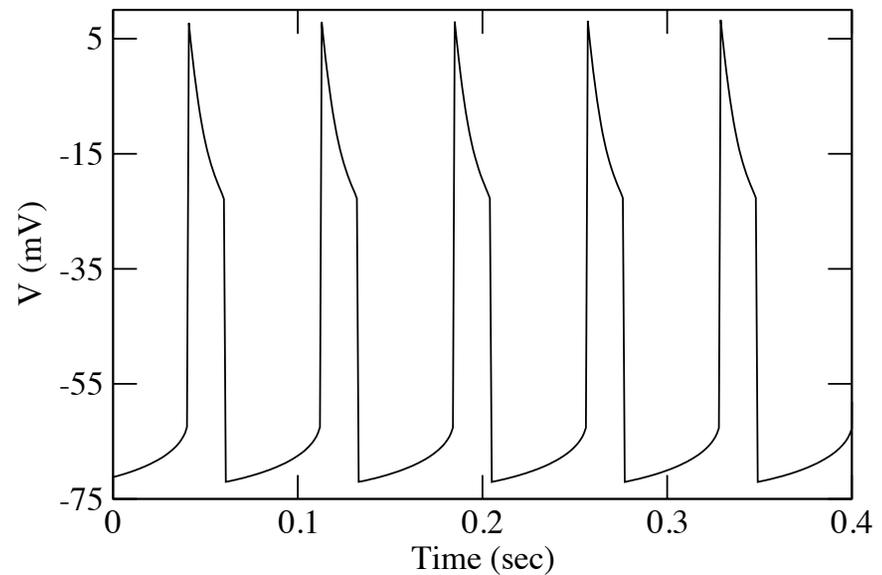
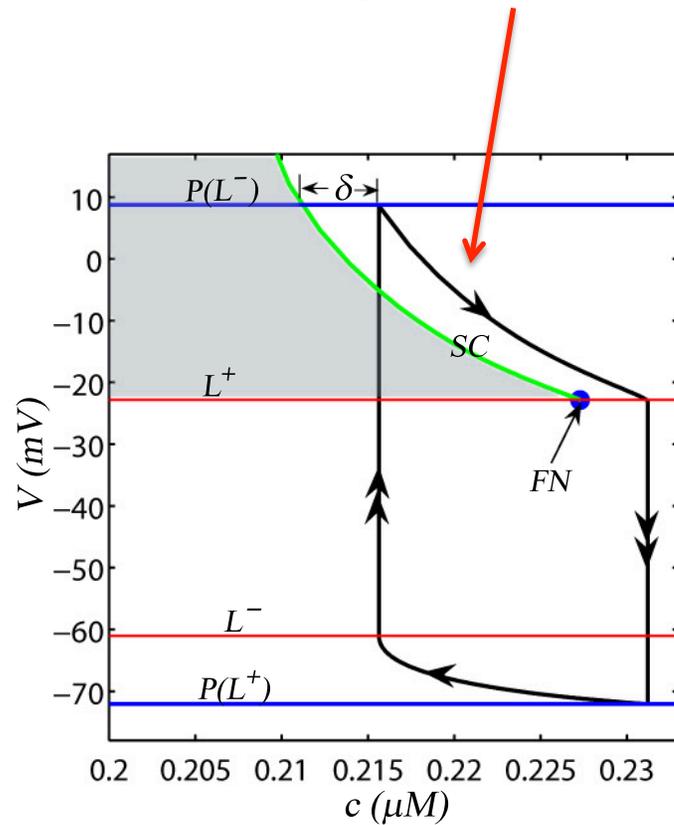
An entire sector of singular canards enter the folded node (FN) from the top (attracting) sheet and travel for some distance along middle (repelling) sheet.

This sector is the **Singular Funnel**, delimited by the fold curve  $L^+$  and the **Strong Canard (SC)**.



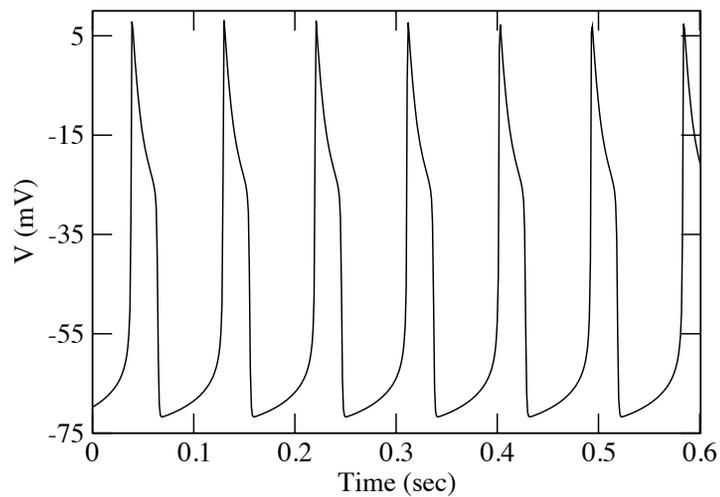
# Relaxation Oscillations

These are periodic solutions that do not enter the singular funnel.

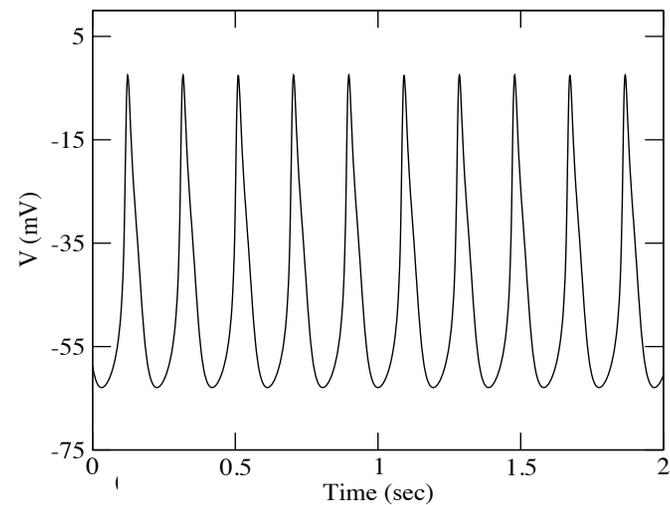


# Continuous Spiking

For  $\varepsilon$  away from 0, the relaxation oscillations transform into a continuous train of impulses.



$C_m = 0.5$  pF

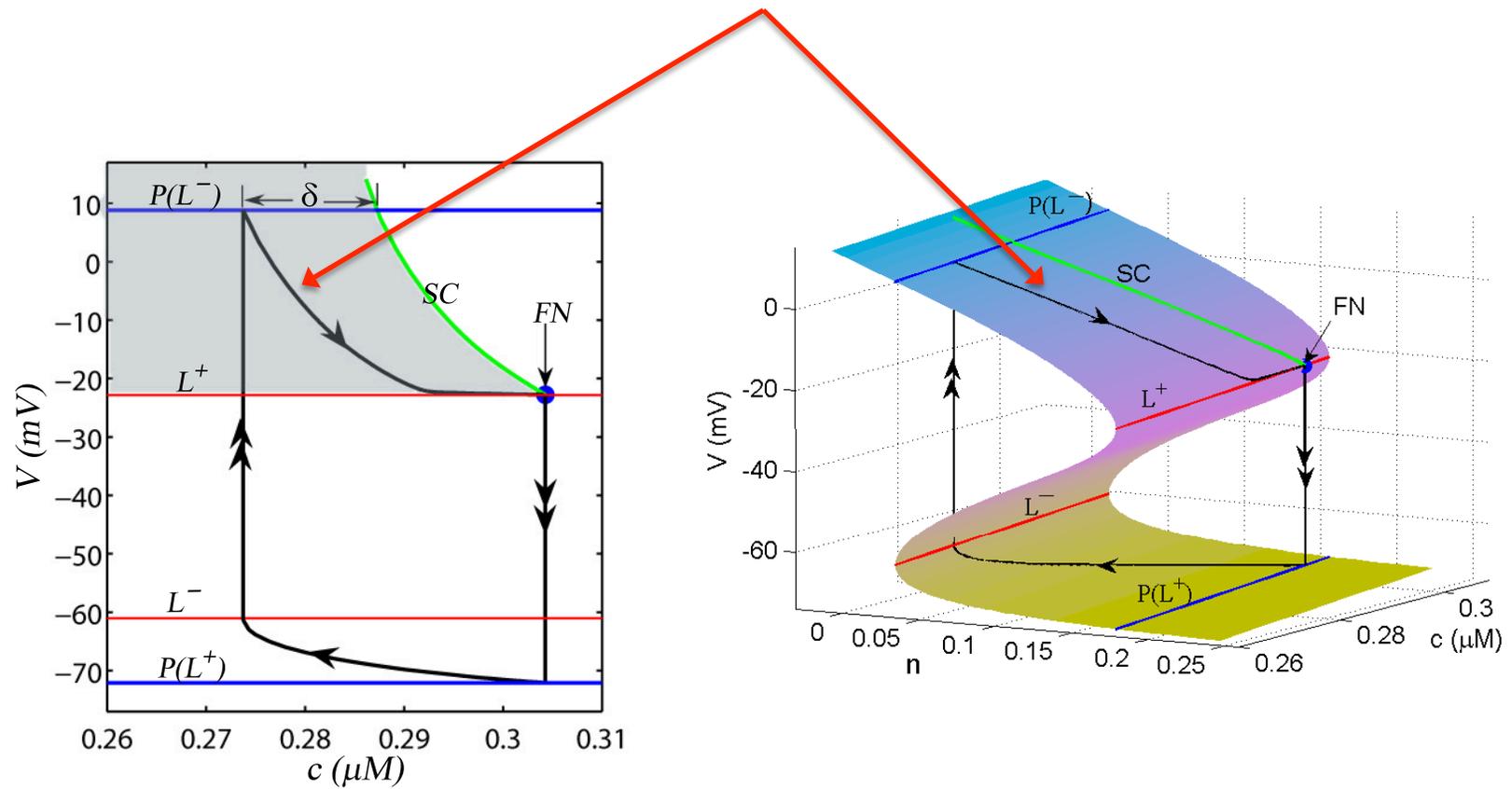


$C_m = 10$  pF

$C_m \approx 5$  pF in lactotrophs/somatotrophs

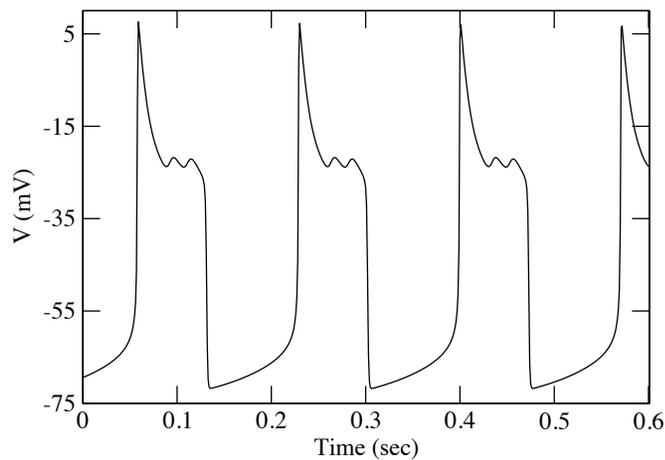
# Mixed Mode Oscillations

These are formed from periodic orbits that enter the singular funnel.

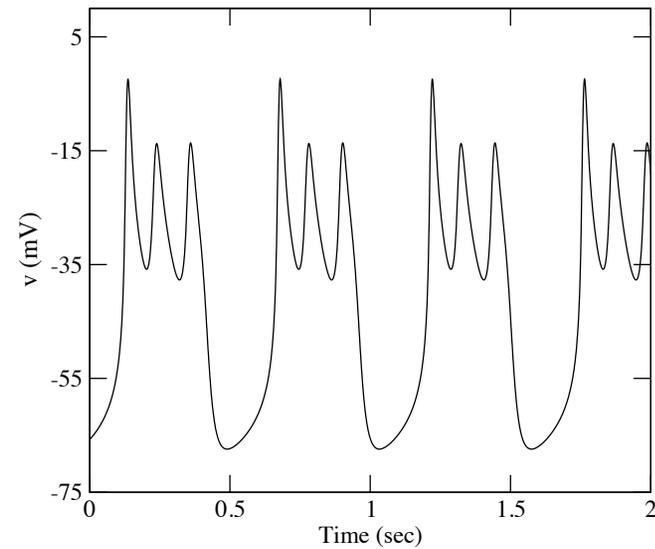


# Pseudo-Plateau Bursting

For  $\epsilon$  away from 0, small oscillations emerge in the vicinity of the folded node. These, combined with the large jumps, form mixed mode oscillations, which in this context, are called **pseudo-plateau bursting**.



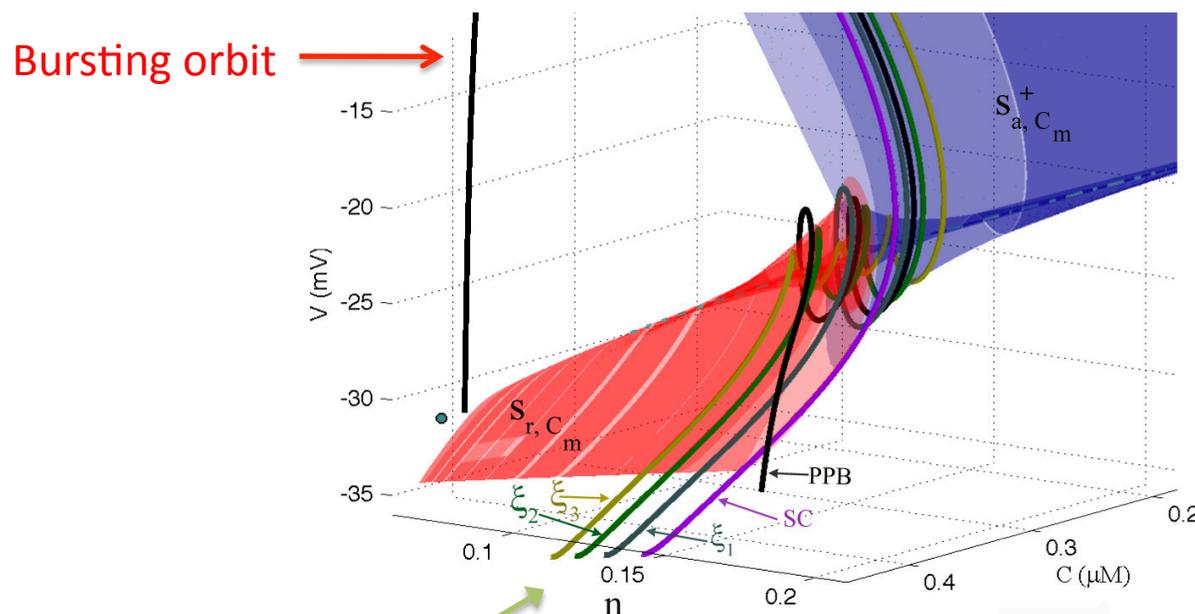
$C_m = 0.5$  pF



$C_m = 10$  pF

# Oscillations Emerge Due to a Twisted Slow Manifold

The sheets of the critical manifold perturb smoothly to form the **slow manifold** for  $\varepsilon > 0$  (Fenichel theory). This is not true in the neighborhood of the folded node, where the perturbed sheets become twisted to preserve uniqueness of solutions.

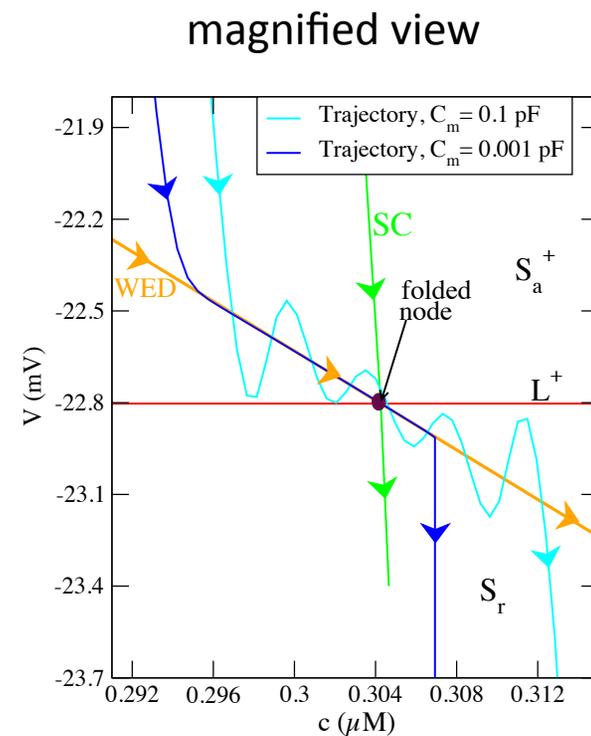
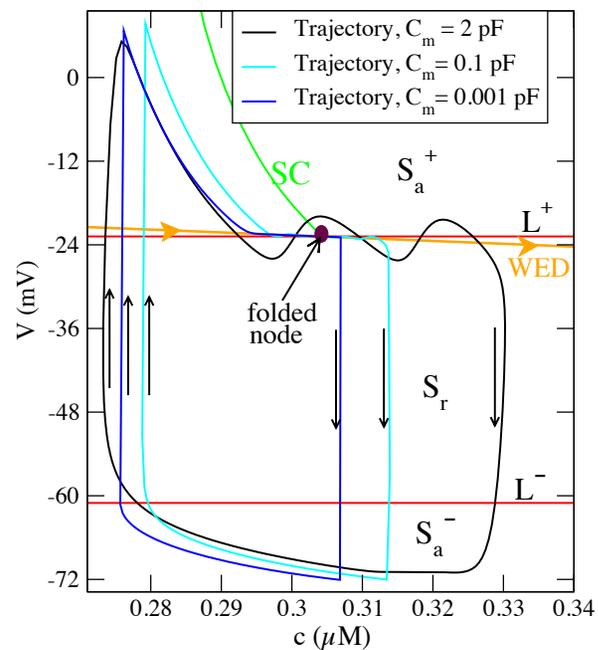


Secondary  
canards

See Desroches et al., Chaos, 18:015107, 2008

# From Singular Orbit to Bursting

Transformation of the periodic orbit as  $\varepsilon$  (or the membrane capacitance  $C_m$ ) is increased.



# Conclusion

The pseudo-plateau bursting oscillations produced by at least some models of pituitary cells are **canard-induced mixed mode oscillations**

This work is soon to be submitted as “What is Pseudo-Plateau Bursting?”, Teka, Tabak, Vo, Wechselberger, Bertram

Go to [Wondimu Teka's talk](#) in Mathematical Neuroendocrinology II (Tuesday at 8:55 AM) to hear more about pseudo-plateau bursting