

Quantum Statistical Mechanics

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Simplest example (finite dimensional case)

$H = \mathbb{C}^M$ $A = M_n(\mathbb{C})$ with ~~some~~ repres: on \mathbb{C}^n

A state $\varphi(a) = \frac{\text{Tr}(pa)}{\text{Tr}(p)}$ with density matrix $p \in M_n(\mathbb{C})$ $p > 0$
 (i.e. $p = \gamma^* \gamma$)

H self adjoint matrix on \mathbb{C}^n ; ~~with~~

\Rightarrow time evolution $\sigma_t(a) = e^{ith} a e^{-ith}$ $\forall t \in \mathbb{R}$

Associated "equilibrium states"

$$\varphi_\beta(a) = \frac{\text{Tr}(a e^{-\beta H})}{\text{Tr}(e^{-\beta H})}$$

$\beta \in \mathbb{R}_+^*$ (inverse temperature)

density $p = e^{-\beta H}$

Gibbs states

$$\underline{\text{equilibrium}}: \varphi(\sigma_t(a)) = \varphi(a) \quad \forall t \in \mathbb{R}$$

Another way to see Gibbs states: Solutions of variational problem for energy functional

Generalizing to infinite dimensional case

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- Observables of a quantum statistical mechanical system:
 $a \in A$ C^* -algebra

- Time evolution $\sigma : \mathbb{R} \rightarrow \text{Aut}(A)$

- Representation $A \rightarrow B(H)$ bounded operators on a Hilbert space H

— is a covariant representation if

$\exists H$ self-adjoint operator on H such that

$$\pi(\sigma_t(a)) = e^{ith} \pi(a) e^{-ith} \quad \forall t \in \mathbb{R}$$

— Positive energy representation if Hamiltonian H can be chosen ~~base~~ with $\text{Spec}(H) \subset [0, \infty)$

- Usually assume A unital; if not \leadsto compactifications

simpliest: 1-point compactification

$$A^{\text{comp}} = A \oplus \mathbb{C} \text{ w/product:}$$

$$(a, \lambda)(b, \mu) = (ab + \lambda b + \mu a, \lambda\mu) \\ a, b \in A \quad \lambda, \mu \in \mathbb{C}$$

Other compactifications:

Stone-Čech compactification

Def: $I \subset A$ essential ideal if

I closed ideal s.t. if $J \subset A$ any closed ideal
 $I \cap J \neq \emptyset$

Def: (Stone-Čech compactification)

A non-unital C^* -algebra, \exists maximal C^* -algebra

$M(A)$ s.t. $A \subset M(A)$ essential ideal

(unique up to isomorphism

"multiplier algebra"

(idea: $M(A)$ consists of all T s.t. $Ta \in A$ for all $a \in A$)
 T bounded lin. operators and $T(ab) = aT(b)$)

Sometime useful to consider also "unbounded multipliers"

$T: D(T) \rightarrow A$ defined on a dense ideal
 $D(T) \subset A$

with $T(ab) = aT(b)$ $\forall a \in A$ and $b \in D(T)$

States: A unital C^* -algebra

$\varphi: A \rightarrow \mathbb{C}$ continuous linear functional
with normalization $\varphi(1)=1$ and positivity
 $\varphi(a^*a) \geq 0 \quad \forall a \in A$

In the non-unital case:

instead of condition $\varphi(1)=1$ have

$$\|\varphi\|=1 \text{ where}$$

$$\|\varphi\| = \sup_{\substack{a \in A \\ \|a\| \leq 1}} |\varphi(a)|$$

GNS representation associated to a state:

given $\varphi: A \rightarrow B(H)$ a repres

\Rightarrow state $\varphi(a) = \langle \xi, \pi(a)\xi \rangle$ for a given $\xi \in H \quad \|\xi\|=1$

Conversely: given a state $\varphi: A \rightarrow \mathbb{C}$

define

$$\langle a, b \rangle = \varphi(a^* b)$$

induces an inner product on the vector space

$$H = A/I \quad I = \begin{array}{l} \text{left ideal} \\ \{a : \varphi(a^* a) = 0\} \end{array}$$

A acts on H by left multiplication

$$a \cdot (b \bmod I) = ab \bmod I$$

$$\{ = 1 + I \quad 1 \in A \text{ unit}$$

$$\varphi(a) = \langle \xi, a\xi \rangle$$

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Equilibrium states for (α, σ_t)

β (inverse temperature)

$$I_\beta = \{ z \in \mathbb{C} \mid 0 < \operatorname{Im}(z) < \beta \}$$

$$\partial I_\beta = \{ z \in \mathbb{C} \mid \operatorname{Im}(z) = 0 \text{ or } \operatorname{Im}(z) = \beta \}$$

$\varphi: A \rightarrow \mathbb{C}$ satisfies the KMS condition
(Kubo-Martin-Schwinger)

at inverse temperature β if

$\forall a, b \in A \quad \exists F_{a,b}(z)$ holomorphic function
on I_β
extending continuously to ∂I_β

such that

$$F_{a,b}(t) = \varphi(a \sigma_t(b))$$

$$F_{a,b}(t+i\beta) = \varphi(\sigma_t(b)a)$$

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Observations:

- $\varphi(\sigma_t(a)) = \varphi(a)$ equilibrium states
- topology on set of KMS_β states Σ_β
 $\varphi_n \rightarrow \varphi$ iff $\varphi_n(a) \rightarrow \varphi(a) \quad \forall a \in A$
 (weak convergence)
 - $\Rightarrow \Sigma_\beta$ compact convex set
 - $\Rightarrow E_\beta \subset \Sigma_\beta$ set of extremal points
 (extremal KMS states)
- A commutative, $\sigma_t = id$ trivial
 \Rightarrow KMS states measures on X $C(X) = A$
 extremal: measures supported on points
- Gibbs states $\varphi(a) = \frac{\text{Tr}(a e^{-\beta H})}{\text{Tr}(e^{-\beta H})}$ are KMS_β -states
 but NOT all KMS_β states are of this form
- $\beta=0$ case : traces i.e.
 $\varphi(ab) = \varphi(ba)$

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Zero temperature KMS - states

Two different notions

Stronger: $\varphi: A \rightarrow C$ KMS $_{\infty}$ if

$\exists \varphi_B$ ~~•~~ $\beta < \infty$ s.t.
(large)

$$\varphi(a) = \lim_{\beta \rightarrow \infty} \varphi_B(a) \quad \forall a \in A$$

Weak limits of positive temperature KMS states

Weaker: (ground states) $\varphi: A \rightarrow C$

$\forall a, b \in A \quad \exists F_{a,b}(z)$ holomorphic on

$$I_B = \{z \in \mathbb{C} \mid \operatorname{Im}(z) > 0\}$$

such that $F_{a,b}(t) = \varphi(a \sigma_t(b))$

too coarse (too many ground states as opposed to KMS $_{\infty}$ -states) not nice convex etc.

- Prop: • $\varphi \in \text{KMS}_{\beta}$ extremal iff GNS rep. is a factor (trivial center)
- Σ_{β} convex compact Choquet simplex w/ extremal points ε_{β}
- Σ_{∞} also convex compact (not a simplex in general)

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Extension of states to multiplier algebra:

Given (A, σ_t) define $M(A)_\sigma \subset M(A)$

multipliers such that $t \mapsto \sigma_t(m)$ norm continuous
 $m \in M(A)_\sigma$

$\varphi: A \rightarrow \mathbb{C}$ state then admits an extension

$$\varphi: M(A)_\sigma \rightarrow \mathbb{C}$$

If $\varphi \in KMS_\beta$ then its extension is still a
 KMS_β -state on $M(A)_\sigma$

If $\mathcal{C} \subset M(A)_\sigma$ s.t. $A \subset \mathcal{C}$ and
 \mathcal{C} is σ_t -invariant then

then $K_p(\mathcal{C}) = KMS_\beta(\mathcal{C})$ can be thought of
as a compactification of $KMS_\beta(A)$

Note: not all $\varphi \in KMS_\beta(\mathcal{C})$ restrict to
elements in $KMS_\beta(A)$ on A : restrict to
positive linear functionals of
 $\|\varphi\| \leq 1$ (not nec. $= 1$)

"quasi-states"

Symmetries:

$$\alpha: G \rightarrow \text{Aut}(A, \sigma_t)$$

$$\alpha \in \text{Aut}(A) \text{ s.t. } \alpha \sigma_t = \sigma_t \alpha \quad \forall t \in \mathbb{R}$$

induced action on KMS states

$$g^* \varphi(a) = \varphi(g(a))$$

$\text{Inn}(A, \sigma_t)$ inner automorphisms

$u \in U(A)$ unitaries $u \in A$ $u^* u = u u^* = 1$

such that $\sigma_t(u) = u \quad \forall t \in \mathbb{R}$

Acting on A by

$$\text{ad}(u) a = u a u^*$$

induced action of $\text{Inn}(A, \sigma_t)$ on KMS states
is trivial:

$$\varphi \in \text{KMS}_\beta \text{ satisfies } \varphi(b \sigma_{ip}(a)) = \varphi(ab) \quad \circledast$$

$\forall a, b$ in a norm-dense σ -invariant subalgebra ("analytic elements")

\circledast is equivalent to KMS condition

$$\text{then } \varphi(u a u^*) = \varphi(a u^* \sigma_{ip}(u)) = \varphi(a)$$

Symmetries by endomorphisms

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(“superselection sectors”)

$\phi : A \rightarrow A$ *-homomorphism
such that

$$\rho^{\sigma_t} = \sigma_t \rho$$

$\rho(1) \neq 1$ in general $\rho(i) = e$ idempotent
such that $\sigma_t(e) = e$

Given $\varphi \in \text{KMS}_\beta$ such that $\varphi(e) \neq 0$

$$\rho^*(\varphi)(a) := \frac{\varphi(\rho(a))}{\varphi(e)} \quad \leftarrow \text{divide to restore normalization}$$

Action of endomorphisms on KMS states

Inner : $\text{Inn}(A, \sigma_t)$ endomorphisms case

~~exists~~ $u \in A$ such that $u^*u = e$
 $u^*u = 1$ (but not uu^* in general now)

and $\sigma_t(u) = \lambda^t u$ for some $\lambda \in \mathbb{R}_+^*$

isometries eigenvectors of time evolution

$$\text{ad}(u)(a) = uau^* \quad \forall a \in A$$

Pf: again $\text{ad}(u)^*(\varphi) = \varphi$ act trivially on KMS states
 $\varphi(uu^*) = \lambda^{-\beta} > 0$ so $\text{ad}(u)^*(\varphi)$ well defined
and then use KMS_β condition as in automorphism case.