

Input also the μ -dependence geometrically

①

A principal $\mathbb{C}^* = \mathbb{G}_m(\mathbb{C})$ -bundle B

$$\mathbb{G}_m \rightarrow B \rightarrow \Delta \quad \text{over the infinitesimal disk } \Delta$$

Note: trivial bundle, but do not fix a choice of a trivialization $\sigma: \Delta \rightarrow B$

Then take $P = B \times \mathbb{G}$ trivial (and trivialized) principal \mathbb{G} -bundle over B

Physical meaning of these data:

\mathbb{G} = group dual to CK Hopf algebra of Feynman graphs

$z \in \Delta$ = complexified dimension of Dim Reg

\mathbb{C}^* = fibre over $z \in \Delta$ is $\mu^{2\epsilon}$

Choice of a unit of mass μ = choice of a section $\sigma: \Delta \rightarrow B$ of fibrations

value $\sigma(\epsilon) = \hbar$

\mathbb{G}_m -action on B by $b \mapsto u(b)$ in a trivialization

$$b = \sigma(z, \sigma(z)) \quad u(b) = (z, u(\sigma(z)))$$

on $P = B \times \mathbb{G}$

$$u(b, g) = (u(b), u^Y(g))$$

Equisingular connections:

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Def: A flat connection on P^0 is equisingular

- iff
- G_m -invariant: $\omega(z, u(v)) = u^Y(\omega(z, v)) \quad \forall u \in G_m$
 - Different sections $\sigma_1, \sigma_2: \Delta \rightarrow B$ with same $\sigma_1(0) = \sigma_2(0)$ have $\sigma_1^*(\gamma) \sim \sigma_2^*(\gamma)$ for γ solution of $D(\gamma) = \omega$

Notation: $V = \pi^{-1}(0) \subset B$ fiber over $z=0$

$B^0 = B \setminus V$ complement

$P^0 = B^0 \times G$

The equivalence relation $\sigma_1^*(\gamma) \sim \sigma_2^*(\gamma)$ is the one of having the same negative piece of the Birkhoff factorization as before

i.e. $\sigma_1^*(\gamma)$ and $\sigma_2^*(\gamma)$ have the same type of singularity at $z=0$ (equisingular)

The first condition

$\omega(z, u(v)) = u^Y(\omega(z, v))$ is a rephrasing

of the scaling property of loops

$$\gamma_{et\mu}(z) = \theta_{tz}(\gamma_\mu(z))$$

while the second condition

$$\sigma_1^*(\gamma) \sim \sigma_2^*(\gamma)$$

is a geometric reformulation of the condition

$$\frac{\partial}{\partial \mu} \gamma_\mu(z) = 0$$

Equivalence relation on equisingular connections (3)

$\omega \sim \omega'$ on P° iff

$$\omega' = Dh + h^{-1}\omega h$$

where h is a G_m -invariant map regular in B

$$h: B \rightarrow G$$

Thm (CM): There is a bijective correspondence between equivalence classes of flat equisingular G -connections on P° and elements $\beta \in \mathfrak{g}(\mathbb{C}) = \text{Lie } G(\mathbb{C})$ with

(1) $\omega \sim D\gamma$

$$(*) \quad \gamma(z, v) = T e^{-\frac{1}{2} \int_0^v u(\beta) \frac{du}{u}}$$

\uparrow
 fixing a trivialization
 $\sigma: \Delta \rightarrow B$

\uparrow
 integration along straight path
 $u = tv \quad t \in [0, 1]$

(2) \circledast Independent of choice of trivialization $\sigma: \Delta \rightarrow B$ with given $\sigma(0)$.

Pf: First show vanishing of two monodromy representations $M_{\Delta^*}(\omega)$ and $M_{\mathbb{C}^*}(\omega)$ around the two generators of $\pi_1(B^\circ)$

(This ensures existence of solutions to $D(x) = \omega$)

Δ^* infinitesimal disk $\Delta^* = \bigcap_n \Delta_n^*$ corresp. to G_n

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- Monodromy $M_{\Delta^* \times \{0\} \times \mathbb{C}^*}(\omega)$ is trivial
because ω extends to $\{z_0\} \times \mathbb{C}$ (only over $\{z=0\}$ singular)

- Monodromy $M_{\Delta^* \times \{u\}}(\omega)$ also trivial
in fact can pick $u=0$ but there
restriction of ω to $\Delta^* \times \{0\}$ is identically zero
(also indep of where choose $u \Rightarrow$ trivial for all u)

\Rightarrow equation $D\gamma = \omega$ has solutions

Notice:

$$\omega(z, u) = u^\gamma(a(z)) dz + u^\gamma(b(z)) \frac{du}{u}$$

$$a, b \in \mathcal{G}(K) \quad \text{flat: } \frac{db}{dz} - \gamma(a) + [a, b] = 0$$

positivity of grading $\Rightarrow \omega$ extends to $\Delta^* \times \mathbb{C}$

Given $(z_0, 0) \in \Delta_n^* \times \{0\}$ construct solution using
path from $(z_0, 0)$ to $(z, 0)$ and then
path $\gamma(z, tv) \quad t \in [0, 1]$

$$\gamma = \gamma_1 \cup \gamma_2$$

$$\gamma^* \omega = \gamma_1^* \omega + \gamma_2^* \omega$$

$$\gamma(z, v) = T e^{\int_0^v u^\gamma(b(z)) \frac{du}{u}}$$

satisfies $\gamma(z, wv) = w^\gamma(\gamma(z, v)) \quad w \in \mathbb{C}^*$

$$\gamma(z, u) = u^\gamma \gamma(z) \quad \gamma(z) = \gamma(z, v) \Big|_{v=1}$$

Since $D\gamma = \omega \Rightarrow \gamma(z)^T d\gamma(z) = a(z) dz, \quad \gamma(z)^T \gamma(z) = b(z)$

Section $\sigma_s(z) = (z, e^{sz}) \quad z \in \Delta$

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$$\sigma_s^*(\gamma)(z) =: \gamma_s(z) = \theta_{sz} \gamma(z)$$

$$\text{and } \frac{\partial}{\partial s} \gamma_s(z) = 0$$

i.e. get loop $\gamma_s(z)$ $\mu = e^s$ in the right class of "physical" loops

then know from previous results $\exists \beta \in \text{Lie}(G(\mathbb{C}))$

s.t.

$$\gamma(z, 1) = T e^{-\frac{1}{2} \int_{\infty}^0 \theta_{-t}(\beta) dt} \cdot \gamma_{\text{reg}}(z)$$

so that

$$\gamma(z, \mu) = v^Y \left(T e^{-\frac{1}{2} \int_{\infty}^0 \theta_{-t}(\beta) dt} \right) v^Y(\gamma_{\text{reg}}(z))$$

$$v^Y \left(T e^{-\frac{1}{2} \int_{\infty}^0 \theta_{-t}(\beta) dt} \right) = T e^{-\frac{1}{2} \int_0^v u^Y(\beta) \frac{du}{u}}$$

$$\Rightarrow \gamma(z, \mu) = \left(T e^{-\frac{1}{2} \int_0^v u^Y(\beta) \frac{du}{u}} \right) v^Y(\gamma_{\text{reg}}(z))$$

use this as gauge equivalence h

$$\Rightarrow \omega \sim D \left(T e^{-\frac{1}{2} \int_0^v u^Y(\beta) \frac{du}{u}} \right)$$

• if $\omega' \sim \omega \quad \omega' = D_h + h^{-1} \omega h \quad \gamma_2(z, \mu) = \sigma_1(z, \mu) h(z, \mu)$

$\Rightarrow \gamma_i$ same reg. part of Birkhoff $\Rightarrow \beta_2 = \beta_1$ (from equality of residues at zero)

• If have $\gamma(z, v) = T e^{-\frac{1}{2} \int_0^v u^\gamma(\beta) \frac{du}{u}}$ ⑥

then $w = D\gamma$ is equisingular (enough check on one solution: others differ by $g \in G(\mathbb{C})$)

Given a section

$\sigma : \Delta \rightarrow B$ $v(z)\sigma(z)$ varying section

$$\gamma_v(z) = T e^{-\frac{1}{2} \int_0^{v(z)} u^\gamma(\beta) \frac{du}{u}}$$

$$= \left(T e^{-\frac{1}{2} \int_0^1 u^\gamma(\beta) \frac{du}{u}} \right) \left(T e^{-\frac{1}{2} \int_1^{v(z)} u^\gamma(\beta) \frac{du}{u}} \right)$$

poly z^{-1}
when paired
w/ $x \in \mathbb{H}$

regular at $z=0$

This shows indep. of choice of trivialization of \otimes

The universal Hopf algebra and group $U^* = U \rtimes G_m$

principal bundles

$$P = B \times G$$

vector bundles

$$E = B \times V$$

$$\rho : G \rightarrow GL(V)$$

Can accommodate simultaneously different G 's (for different theories)

Def: (E, W) filtered vector bundle E on B
increasing filtration

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$$W^{-n-1}(E) \subset W^{-n}(E)$$

and trivialization of associated graded

$$\text{Gr}_n^W(E) = W^{-n}(E) / W^{-n-1}(E)$$

A W -connection ∇ on E :

connection on $E^\circ = E|_{B^\circ}$ (i.e. singular on $\pi^{-1}(0)$)

Such that

- 1) ∇ compatible w/ filtration: restricts to a connection on each $W^{-n}(E^\circ)$
- 2) ∇ induces the trivial connection on $\text{Gr}_n^W(E)$

Equivalence: $\nabla_1 \sim \nabla_2$ on E° (W -equivalence)

iff $\exists T \in \text{Aut}(E)$ s.t. T preserves the filtration, induces identity on $\text{Gr}_n^W(E)$, and

$$T \circ \nabla_1 = \nabla_2 \circ T$$

intertwines the connections

Equisingular : ∇

$V = \text{fin. dim. } \mathbb{Z}\text{-graded vector space}$

$$V = \bigoplus_{n \in \mathbb{Z}} V_n$$

$E = B \times V \quad (E, W)$ with

$$W^{-n}(E) = \bigoplus_{m \geq n} V_m$$

G_m -action induced by grading

Def: ∇ a W -connection on $E = B \times V$ equisingular
iff G_m -invariant and

$\nabla \eta = 0$ solutions (fundam. system of solutions
= basis of flat sections)

$$\sigma_1^*(\eta) \sim \sigma_2^*(\eta) \quad \text{for } \sigma_1(0) = \sigma_2(0)$$

$\sigma_i : \Delta \rightarrow B$

A category of flat equisingular
vector bundles (E, ∇)

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\mathcal{E} Objects of \mathcal{E} : pairs $\Theta = [V, \nabla]$
W-equivalence class
of ∇

Morphisms $\text{Hom}_{\mathcal{E}}(\Theta, \Theta')$

$T: V \rightarrow V'$ compatible w/ grading and st
on the bundle $(E' \oplus E)^*$ the W-connection

$$\nabla_1 := \begin{pmatrix} \nabla' & 0 \\ 0 & \nabla \end{pmatrix} \text{ and}$$

$$\begin{aligned} \nabla_2 &= \begin{pmatrix} 1 & T \\ 0 & 1 \end{pmatrix}^{-1} \nabla_1 \begin{pmatrix} 1 & T \\ 0 & 1 \end{pmatrix} = \\ &= \begin{pmatrix} \nabla' & T\nabla - \nabla'T \\ 0 & \nabla \end{pmatrix} \end{aligned}$$

are W-equivalent: $\nabla_2 \sim \nabla_1$

To avoid a problem with morphisms of filtered
vector spaces (not abelian category)

- If G dual to $\mathcal{H} = \bigoplus_{n \geq 0} \mathcal{H}_n$
 ω flat equivariant on $P = B \times G$

\Rightarrow given $\xi: G \rightarrow GL(V)$ repres.

(ω, ξ) determine an element $\Theta = [E, \nabla] \in \text{Obj}(\mathcal{E})$

and $\omega \sim \omega' \Rightarrow$ same Θ

Notice: an affine group scheme G is completely determined by its category of finite dimensional linear representations \mathcal{R}_G

So replacing $P = B \times G$ by all the possible $E = B \times V$ with $\xi: G \rightarrow GL(V)$ maintains same information.

\Rightarrow way to recover each G for each particular theory inside a common \mathcal{E} category

Theorem: $\mathcal{E} \cong \mathcal{R}_{U^*}$ finite dim linear representations
 (CM) $U^* = U \rtimes G_m$ where

$L_U = F(1, 2, 3, \dots)_\bullet$ = free graded Lie algebra with one generator e_{-n} in each degree $n \geq 1$

$H_U = U(F(1, 2, 3, \dots)_\bullet)^\vee$ Hopf algebra

Renormalization group

$$e = \sum_{n=1}^{\infty} e_n \quad (\text{nilpotent Lie algebra})$$

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$$\text{rg}: G_a \rightarrow \mathbb{D} \quad \text{additive subgroup generated by } e$$

$$\text{Beta function} \quad \beta = \sum_{n=1}^{\infty} \beta_n \quad \in \text{Lie } G$$

components of $\text{deg} = n$ of beta function

$$\Rightarrow e_n \mapsto \beta_n$$

morphism from

$$\text{Lie } \mathbb{D} \rightarrow \text{Lie } G$$

integrates to morphism $\mathbb{D} \rightarrow G$

then \mathbb{D} maps to all G 's of different physical theories through their β -function

"Universal counterterms"

