

# The spherical ensemble with external sources

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# 1 Outline

- ① Introduction
- ② Potential theory
- ③ Vector equilibrium problem
- ④ About the proof

# 1 Introduction: The spherical ensemble

$n$  random points on the unit sphere  $\mathbb{S}^2 \subset \mathbb{R}^3$  with joint probability density

$$\frac{1}{Z_n} \prod_{i=1}^{n-1} \prod_{j=i+1}^n \|x_i - x_j\|^2$$

- ▶ Determinantal point process
- ▶ Eigenvalues of a random matrix Krishnapur (2009)

Take independent Ginibre matrices  $A$  and  $B$ .

Compute eigenvalues of  $AB^{-1}$ .

Put points on  $\mathbb{S}^2$  with inverse stereographic projection.

# 1 The spherical ensemble with external charges

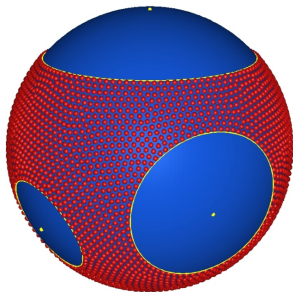
Fix  $p_0, \dots, p_r \in \mathbb{S}^2$  and  $a_0, \dots, a_r > 0$ . It creates probability density for points  $x_1, \dots, x_n$  on the sphere

$$\frac{1}{Z_n} \prod_{i=1}^{n-1} \prod_{j=i+1}^n \|x_i - x_j\|^2 \cdot \prod_{j=1}^n \prod_{k=0}^r \|x_j - p_k\|^{2na_k}$$

Each  $p_k$  creates a region around itself that is free from  $x_j$ 's in large  $n$  limit.

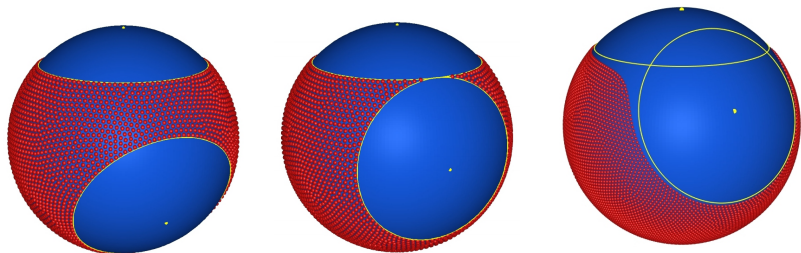
When  $a_k$ 's are small these regions are spherical caps

Brauchart, Dragnev, Saff,  
Womersley



## 1 The spherical ensemble with external charges

When charges  $a_k$  are larger the regions start to overlap.



Figures from BDSW

- ▶ What region  $D$  is occupied by the free charges?
- ▶  $D$  is known as the **droplet**.

## 1 Normal matrix model

Eigenvalues of **complex Ginibre matrix** have joint density

$$\frac{1}{Z_n} \prod_{i=1}^n \prod_{j=i+1}^n |z_i - z_j|^2 \prod_{j=1}^n e^{-n|z_j|^2}$$

- Eigenvalues fill out a disk as  $n \rightarrow \infty$ .

Ginibre model with **external sources** (also arises as eigenvalues in a normal matrix model)

$$\frac{1}{Z_n} \prod_{i=1}^n \prod_{j=i+1}^n |z_i - z_j|^2 \prod_{j=1}^n e^{-n|z_j|^2} \prod_{j=1}^n \prod_{k=0}^r |z_j - q_k|^{2na_k}$$

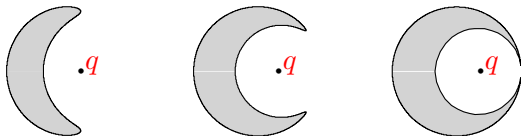
- Studied in detail for case  $r = 0$  by

Balogh, Bertola, Lee, McLaughlin 2015

# 1 Droplet

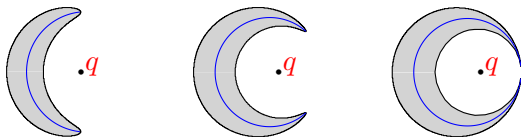
$$\frac{1}{Z_n} \prod_{i=1}^n \prod_{j=i+1}^n |z_i - z_j|^2 \prod_{j=1}^n e^{-n|z_j|^2} \prod_{j=1}^n |z_j - q|^{2na}$$

Change in **droplet** as  $a$  increases.



# 1 Motherbody

**Average characteristic polynomial**  $P_n(z) = \mathbb{E} \left[ \prod_{j=1}^n (z - z_j) \right]$   
is **planar orthogonal polynomial** with zeros that cluster around a contour as  $n \rightarrow \infty$ , known as the **motherbody**



- Planar orthogonality can be rewritten as orthogonality along a closed contour around  $q$ .

Balogh, Bertola, Lee, McLaughlin 2015, Lee, Yang 2019

- This allows a Riemann-Hilbert asymptotic analysis



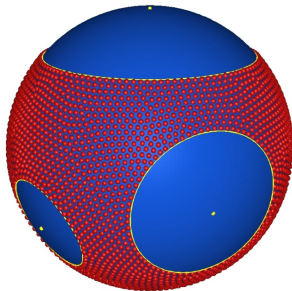
## 2 Outline

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## Potential theory

### Back to spherical model

$$\frac{1}{Z_n} \prod_{i=1}^{n-1} \prod_{j=i+1}^n \|x_i - x_j\|^2 \cdot \prod_{j=1}^n \prod_{k=0}^r \|x_j - p_k\|^{2i}$$



**Fixed charge distribution**  $\sigma = \sum_{k=0}^r a_k \delta_{p_k}$

**Limiting distribution**  $\mu_\sigma$  of free charges is unique probability measure on  $\mathbb{S}^2$  that satisfies

$$\begin{aligned} U^{\mu_\sigma} + U^\sigma &= \ell & \text{on } D = \text{supp}(\mu_\sigma), \\ U^{\mu_\sigma} + U^\sigma &\geq \ell & \text{on } \mathbb{S}^2. \end{aligned}$$

**Notation:**  $U^\mu(x) = \int \log \frac{1}{\|x - y\|} d\mu(y)$

## 2 Potential theory

**Fixed charge distribution**  $\sigma = \sum_{k=0}^r a_k \delta_{p_k}$

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Lemma

$$\mu_\sigma = (\lambda(D))^{-1} \lambda_D$$

where  $\lambda_D$  is the restriction to  $D$  of the normalized Lebesgue measure  $\lambda$  on  $\mathbb{S}^2$ .

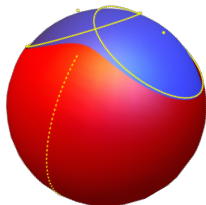
## 2 The motherbody

A **motherbody** for  $D$  is a probability measure  $\sigma^*$  on a one-dimensional subset of  $D$  such that

$$\begin{aligned} U^{\sigma^*} &= (\lambda(D))^{-1} U^{\lambda_D} + \ell^* && \text{on } \mathbb{S}^2 \setminus D, \\ U^{\sigma^*} &\geq (\lambda(D))^{-1} U^{\lambda_D} + \ell^* && \text{on } D. \end{aligned}$$

In case of two points  $p_0, p_1$  with sufficiently large equal charges  $a_0 = a_1$ ,  $\mathbb{S}^2 \setminus D$  is an **ellipse** after stereographic projection onto the complex plane.

The motherbody is minimizer of **equilibrium problem with external field**  
**Criado del Rey - K**



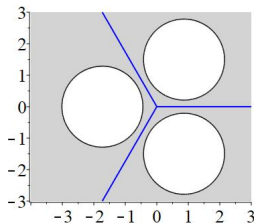
## 2 More external charges

We generalize this to  $r + 1$  points  $p_0, \dots, p_r$  **with symmetry**

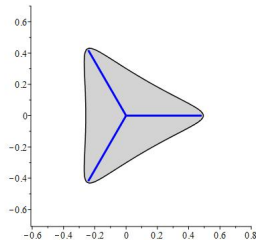
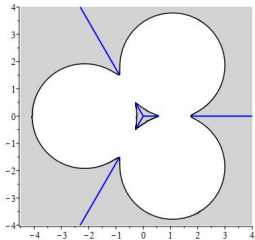
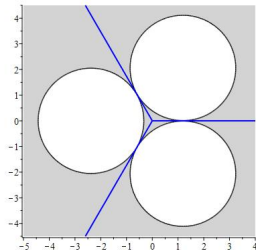
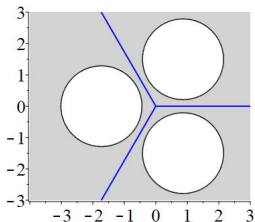
- ▶  $p_j$ 's are located symmetrically around the north pole
- ▶ Equal charges  $a_0 = a_1 = \dots = a_r = a > 0$

Motherbody will be supported on  $r + 1$  **meridians** on the sphere, connecting the north and south poles.

After stereographic projection these are  $r + 1$  **half-rays** in the complex plane.



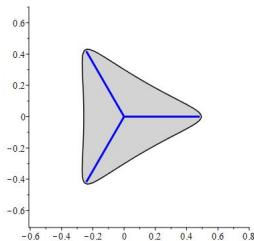
## 2 More pictures for $r = 2$ (three charges)



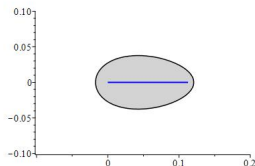
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### 3 Remove the symmetry



$$z \mapsto z^{r+1}$$



We expect **three cases** for the support  $\Sigma$  of the motherbody (after removing the symmetry)

- ▶ **Bounded interval support (BIS):**  $\Sigma = [0, x_1]$
- ▶ **Two interval support (TIS):**  $\Sigma = [0, x_1] \cup [x_2, \infty)$
- ▶ **Full interval support (FIS):**  $\Sigma = [0, \infty)$



### 3 Vector equilibrium problem with $r$ measures

**Energy functional** with input  $q > 0$ ,  $0 < t < 1$  and integer  $r \geq 1$

$$\begin{aligned}\mathcal{E}(\mu_1, \dots, \mu_r) &= \sum_{j=1}^r I(\mu_j) - \sum_{j=1}^{r-1} I(\mu_j, \mu_{j+1}) \\ &\quad - \frac{1-t}{t} \int \log |x + q^{-1}| d\mu_1(x) + \frac{r+t}{t} \int \log |x - (-1)^r q| d\mu_r(x)\end{aligned}$$

**Notation:**

$$\begin{aligned}I(\mu, \nu) &= \iint \log \frac{1}{|x - y|} d\mu(x) d\nu(y) \\ I(\mu) &= I(\mu, \mu)\end{aligned}$$

### 3 Vector equilibrium problem with $r$ measures

**Energy functional** with input  $q > 0$ ,  $0 < t < 1$  and integer  $r \geq 1$

$$\begin{aligned}\mathcal{E}(\mu_1, \dots, \mu_r) &= \sum_{j=1}^r I(\mu_j) - \sum_{j=1}^{r-1} I(\mu_j, \mu_{j+1}) \\ &\quad - \frac{1-t}{t} \int \log |x + q^{-1}| d\mu_1(x) + \frac{r+t}{t} \int \log |x - (-1)^r q| d\mu_r(x)\end{aligned}$$

**Minimize**  $\mathcal{E}(\mu_1, \dots, \mu_r)$  under conditions

►  $\mu_j$  is supported on  $[0, \infty)$  if  $j$  is odd, and on  $(-\infty, 0]$  if  $j$  is even,

►  $\int d\mu_j = 1 + \frac{j-1}{t}$  for  $j = 1, \dots, r$ .

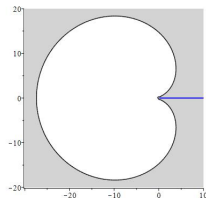
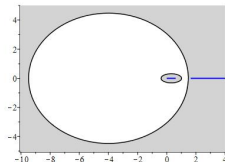
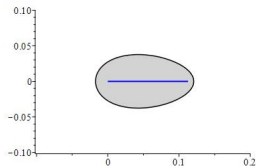
### 3 Main result

#### Theorem

- (a) *The vector equilibrium problem has a **unique minimizer**.*
- (b) *The first component  $\mu_1$  has **support**  $\Sigma_1$  that is either  $[0, x_1]$ , or  $[0, x_1] \cup [x_2, \infty)$  or  $[0, \infty)$ .*
- (c)  *$\Phi(z) = t \int \frac{d\mu_1(s)}{z-s} + \frac{1-t}{z+q^{-1}}$  is an **algebraic function** with meromorphic continuation to an  $r+1$  sheeted **Riemann surface***
- (d) *There is a **domain**  $U$  containing  $\Sigma_1$  with boundary*

$$\partial U : \quad \boxed{z\Phi(z) = \frac{|z|^{\frac{2}{r+1}}}{1 + |z|^{\frac{2}{r+1}}}}$$

### 3 Pictures (on different scales)



**BIS:** closed contour    **TIS:** two contours    **FIS:** closed contour

$\partial U$  is boundary of **shaded region**  $U$

- ▶ **BIS:**  $\Sigma_1 = [0, x_1]$  and  $U$  is bounded
- ▶ **TIS:**  $\Sigma_1 = [0, x_1] \cup [x_2, \infty)$  and  $U$  has one bounded and one unbounded component
- ▶ **FIS:**  $\Sigma_1 = [0, \infty)$  and  $U$  is unbounded

### 3 Main result (cont.)

#### Theorem

(e)  $\Omega = \{z \in \mathbb{C} \mid z^{r+1} \in U\}$  is mapped by inverse stereographic projection to the **droplet**  $D$

The points  $p_0, \dots, p_r$  correspond to **equidistant points** in the plane on a circle with **radius**  $q^{-\frac{1}{r+1}} > 1$

The **charges** are

$$a = \frac{1-t}{t(r+1)} > 0$$

(f) The measure  $\mu_1$  gives the **motherbody** of  $D$  (after re-introducing the  $r+1$ -fold symmetry and mapping back to the unit sphere)

## 4 Outline

- ① Introduction
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## 4 Vector equilibrium problem

- ▶ Vector equilibrium problem is **weakly admissible**
- ▶ The measures  $\mu_2, \dots, \mu_r$  have full supports  $(-\infty, 0]$  or  $[0, \infty)$  as they are **balayage measures**

$$2\mu_k = \text{Bal}(\mu_{k-1} + \mu_{k+1}; (-1)^{k+1}[0, \infty)), \quad \text{for } k = 2, 3, \dots, r$$

## 4 Vector equilibrium problem

- ▶ The measures  $\mu_2, \dots, \mu_r$  have full supports  $(-\infty, 0]$  or  $[0, \infty)$  as they are **balayage measures**
- ▶  $\mu_1$  also minimizes

$$\begin{aligned} & \frac{1}{2} \iint \log \frac{1}{|x - y|} d\mu(x) d\mu(y) \\ & \quad + \frac{1}{2} \iint \log \frac{1}{|x^{1/r} - y^{1/r}|} d\mu(x) d\mu(y) \\ & + \int \left( \left(1 - \frac{1}{t}\right) \log \left(x + q^{-1}\right) + \left(1 + \frac{r}{t}\right) \log \left(x^{1/r} + q^{1/r}\right) \right) d\mu(x) \end{aligned}$$

among probability measures  $\mu$  on  $[0, \infty)$

Equilibrium problem for **Muttalib-Borodin** ensemble

Claeys Romano 2014, K 2016, Molag 2020



## 4 The support of $\mu_1$

$\mu_1$  is the balayage measure  $\text{Bal}\left(-(\frac{1}{t}-1)\delta_{-q} + \mu_2; [0, \infty)\right)$   
provided that this is **positive**

► This happens in FIS (full interval support) case.

There are four possibilities for the **support**  $\Sigma_1$  of  $\mu_1$

►  $\Sigma_1 = [0, x_1]$  with  $0 < x_1 < \infty$

**BIS**

►  $\Sigma_1 = [0, x_1] \cup [x_2, \infty)$  with  $0 < x_1 < x_2 < \infty$

**TIS**

►  $\Sigma_1 = [0, \infty)$

**FIS**

►  $\Sigma_1 = [x_2, \infty)$  with  $0 < x_2 < \infty$

**UIS**

**UIS** (unbounded interval support) does not happen for  $0 < q \leq 1$ , **BIS** does not happen for  $q \geq 1$ .

This is proved with **iterated balayage** algorithm

Dragnev Kuijlaars 1999

## 4 Monotonicity properties

$\mu_1 = \mu_{1,t}$  and  $\Sigma_1 = \Sigma_{1,t}$  vary with  $t \in (0, 1)$

Fix  $0 < q < 1$

### Theorem

- (a)  $\Sigma_{1,t}$  *increases with*  $t$ .
- (b)  $t\mu_{1,t}$  *increases with*  $t$ .
- (c) **There are two critical values**  $0 < t_{1,cr} < t_{2,cr} < 1$  **such that**
  - BIS** in case  $0 < t \leq t_{1,cr}$
  - TIS** in case  $t_{1,cr} < t < t_{2,cr}$
  - FIS** in case  $t_{2,cr} \leq t < 1$

(a) is implied by (b) and (b) follows from further analysis of iterated balayage algorithm (+some other tricks)

(c) follows from explicit calculations in BIS and FIS cases.

## 4 Riemann surface

**Riemann surface  $\mathcal{R}$  with  $r + 1$  sheets**

$$(\Sigma_j = \text{supp}(\mu_j))$$

$$\mathcal{R}^{(1)} = \mathbb{C} \setminus \Sigma_1,$$

$$\mathcal{R}^{(j)} = \mathbb{C} \setminus (\Sigma_{j-1} \cup \Sigma_j) \quad \text{for } j = 2, \dots, r,$$

$$\mathcal{R}^{(r+1)} = \mathbb{C} \setminus \Sigma_r.$$

► **Meromorphic function  $\Phi$  is given on  $j$ th sheet by**

$$\Phi^{(j)}(z) = t \int \frac{d\mu_j(s)}{z-s} - t \int \frac{d\mu_{j-1}(s)}{z-s}, \quad z \in \mathcal{R}^{(j)}$$

**with  $\mu_0 = \left(1 - \frac{1}{t}\right) \delta_{-q^{-1}}$  and  $\mu_{r+1} = \left(1 + \frac{r}{t}\right) \delta_{(-1)^r q}$**

►  **$z\Phi(z)$  has degree 2.**

## 4 Droplet and spherical Schwarz function

- ▶  $\partial U : z\Phi(z) = \frac{|z|^{\frac{2}{r+1}}}{1 + |z|^{\frac{2}{r+1}}}$  is boundary of domain  $U$ .
- ▶ If  $\Omega = \{z \mid z^{r+1} \in U\}$  and  $S(z) = z^r \Phi(z^{r+1})$  then

$$\partial\Omega : S(z) = \frac{\bar{z}}{1 + |z|^2}$$

$S(z)$  is the **spherical Schwarz function** of  $\partial\Omega$

- ▶ **Monotonicity** of  $\Omega$  in parameter  $t$  is crucial to prove that  $\Omega$  is the stereographic projection of the **droplet**.

## 4 Final remark

**Thank you for your attention !**